## Glass transition temperature of poly(methyl methacrylate) blends

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The glass transition temperature  $(T_a)$  of atactic poly(methyl methacrylate) (PMMA) blends was measured over a wide range of molecular weights  $(\bar{M}_n)$  via differential thermal analysis. Comparison of these results critically questions the validity of the Fox-Flory relationship as anything other than a first order approximation of a more complex functional relationship.

In the last twenty-five years, several approaches have been advocated to explain the relationship between molecular size/shape and the second-order transition associated with amorphous polymers<sup>1-6</sup>. By far the two most popular theories have been the free-volume theory of Fox and Flory<sup>3,4</sup> and the statistical mechanical approaches of Gibbs and DiMarzio<sup>5,6</sup>. Various degrees of success have been reported using these theories. The Fox-Flory approach generally is used to approximate  $T_g$  vs  $M_n^{-1}$  either for  $M_n > 10^4$  <sup>7.8</sup> or as a series of straight line segments <sup>1,2,9</sup>, and the Gibbs-DiMarzio theories set either the fractional free-volume at  $T_a(V_0)$  or the energy barrier to segmental rotation ( $\epsilon$ ) to a constant (most commonly with  $V_0 = 0.025^{10,11}$ ) and floating the other parameter to obtain the best fit. The present study considers both theories, although the statistical methods vield more consistent results.

Three feedstocks were used to prepare some 24 blends by precipitation in a hexane solution after dissolution in acetone at 37°C (Table 1)<sup>12</sup>. These were designated high (H), medium (M), and low (L) molecular weights. Both H and M were obtained commercially; L was prepared in this laboratory by solution polymerization in toluene at 60°C<sup>13,14</sup>. Molecular weights were determined by g.p.c. in THF using polystyrene standards. Additionally, a PMMA secondary standard was evaluated using the same g.p.c. procedure and compared with the reported light scattering ( $\bar{M}_w = 60.6 \times 10^3$ ) and membrane osmometry results ( $\bar{M}_n = 33.2 \times 10^3$ )<sup>15</sup>. Glass transition temperatures for each vacuum-dried blend were obtained using a Dupont 990 Differential Thermal Analyser (d.t.a.) at a

Table 1 Characterization of PMMA Feedstocks

Designation	$\overline{M}_W \times 10^{-3}$	$\overline{M}_n \times 10^{-3}$	$\overline{M}_W/\overline{M}_D$
H <sup>a</sup>	364	122	2.99
Mb	64.4	32.1	2.01
L	5.03	3.90	1.30
S <sup>c</sup>	61.5 (60.6)	29.6 (33.2)	2.07 (1.82)

<sup>&</sup>lt;sup>a</sup> Aldrich Chemical Company, Milwaukee, Wis. (cat. no. 18, 226–5,

heating rate of 20°C min<sup>-1</sup> under nitrogen (0.2 l min<sup>-1</sup>). Multiple run results for each blend combination (e.g. H/M) are shown in Figure 1. Here regression lines for H/M, H/L, and M/L are highly significant (0.001 level), the  $T_a$ s averaging 395, 382 and 347 K for H, M, and L, respectively.

Based on the previous observations of heterogeneous systems, that the mean number-average molecular weight is the primary consideration, 3,4,6,16,17 a reduction of the  $T_g$  vs composition plots to  $T_g$  vs  $\bar{M}_n^{-1}$  via  $\bar{M}_n^{-1} = (w/\bar{M}_n)_i$  $+(w/M_n)_i$  (where i and j refer to the feedstock components and w are their respective weight fractions) is shown in Figure 2. Once again the Fox-Flory expression yields lines that are significant at the 0.001 level, although the

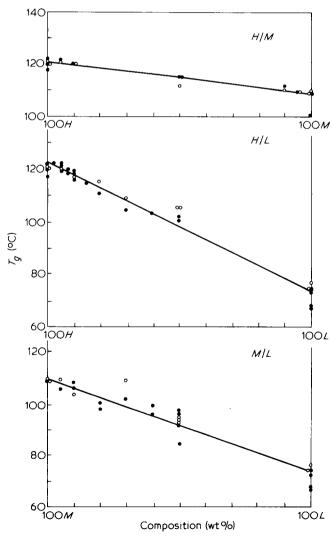


Figure 1 Dependence of  $T_q$  on blend composition:  $\bullet$ ,  $\circ$ , denote the first and second d.t.a. scans, respectively

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<sup>[</sup> $\eta$ ] = 1.3 for  $\overline{M}_W$  = 500 × 10<sup>3</sup>) b Polysciences, Warrington, Pa. (cat. no. 4554, lot no. 1888–125,  $[\eta] = 0.2 \text{ for } \overline{M}_W = 50.0 \times 10^3)$ 

<sup>&</sup>lt;sup>c</sup> Aldrich Chemical Company, Milwaukee, Wis. (cat. no. 18, 225-7, lot no. 03)

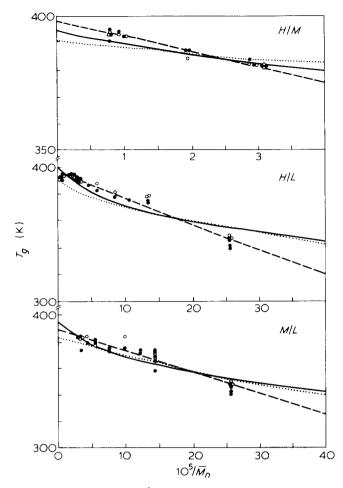


Figure 2 Influence of  $M_n^{-1}$  on  $T_g$  and the associated theoretical analyses: - - -, Fox-Flory; -- $\stackrel{\circ}{-}$ , Gibbs, and  $\cdots$ , Gibbs-DiMarzio (Table 2)

three slopes differ by a factor of > 3 (*Table 2*). In contrast for  $V_0 = 0.025$ , the Gibbs and Gibbs-DiMarzio relationships gives  $\varepsilon = 0.99-1.00$  and  $\varepsilon = 1.66-1.69$  kcal (mol segments)<sup>-1</sup>, respectively. The extrapolated  $T_q$ s for the infinite molecular weight polymer,  $T_q^{\infty}$ , are displayed in parentheses in *Table 2*.

As Figure 2 shows, each set of results appear to fit the Fox-Flory relationship best. However, these seemingly good fits do not yield consistent values for the constants,  $T_a^{\infty}$  and  $K_a$ . Since any theoretical argument should uniquely describe a relationship that is independent of molecular weight (at least for molecular weights greater than the entanglement molecular weight, i.e.  $\bar{M}_n \simeq 10^4$ (Table 2 of ref 9)), the only recourse is to suggest that these apparent straight line relationships really represent approximations of some curvilinear function. This hypothesis is borne out by the observation that both  $K_a$  and  $T_a^{\infty}$  decrease as the  $\bar{M}_n$ s of the respective blend combinations extend into the lower molecular weight regions. Earliest results by Fox and Loshaek <sup>18</sup> predict  $K_g = 0.6 - 1.5 \times 10^5 \text{K}$ , while Pratt<sup>19</sup> reported  $1.55 \times 10^5 \text{K}$  for 7-irradiated PMMA, Beevers and White<sup>20</sup> found 2.1  $\times$  10<sup>5</sup>K for a molecular weight range of 2.7–77  $\times$  10<sup>3</sup>, and Thompson<sup>21</sup> observed 3.0 and  $3.4 \times 10^5$ K for 0.75 and 0.82 s-PMMA, respectively. Interestingly our three basic blends fall within the scatter envelope formed by the compilation of data presented by Boyer (Figure 1, ref 2). Moreover, the present results indicate that, in the highest

Table 2 Theoretical comparison of dependence of  $T_q$  on  $\overline{M}_n$  for

	Fox-Flory	Gibbs <sup>a</sup>	Gibbs—DiMarzio <sup>a</sup>
H/M	$K_g^b = 5.402$ (398.0) $d$	$e^{c} = 0.99$ (395.1)	$\epsilon = 1.69$ (391.2)
H/L	<i>K<sub>g</sub></i> = 1.952 (397.1)	$\epsilon = 1.00$ (399.1)	ε = 1.68 (388.9)
M/L	<i>K<sub>g</sub></i> = 1.588 (388.1)	$\epsilon = 0.99$ (395.1)	$\epsilon = 1.66$ (384.2)

<sup>&</sup>lt;sup>a</sup> Evaluated for  $V_0 = 0.025$ 

temperature region,  $K_g$  may increase at a greater rate than Boyer's plot indicates. The implication, that  $T_q^{\infty}$  is asymptotically approaching a maximum, can be rationalized on the grounds that there must be a maximum permissible segmental chain stiffness above which coordinated chain motion does not occur.

With the foregoing arguments in mind, the Gibbs and Gibbs-DiMarzio functions were fitted (Figure 2) assuming an iso-free-volume state (i.e.  $V_0 = 0.025$ ). Although the data point to the theoretical line correspondence is not as good as the Fox-Flory relationship, a uniform value for  $\varepsilon$ prevailed,  $\Delta \varepsilon$  varying by only 0.01 kcal (mol segment)<sup>-1</sup> between H/M, H/L, M/L (Table 2). Using d.t.a., several authors <sup>17,21</sup> have shown that  $\varepsilon = 0.98-1.00$  kcal (mol segments)<sup>-1</sup>, further confirming the solidarity of the observations. Similar agreement between past and present work is noted for the Gibbs-DiMarzio relationships, wherein  $\varepsilon = 1.66$  kcal (mol segment)<sup>-1</sup> is typical<sup>20</sup>. Further support is found in the observation that both these functions are sensitive to changes in  $\varepsilon$ ,  $V_0$ , or the degree of polymerization.

These experiments question the validity of the freevolume theory of Fox-Flory as the explanation for the  $T_a$ vs  $\bar{M}_n^{-1}$  of PMMA. Although both within the present context and with respect to the general acrylic literature the Gibbs theory yields the most consistent results (cf overview in ref 22\*), that fit is not totally satisfactory either. One approach which this latter observation has fostered is a systematic reexamination of the statistical mechanical relations. While a common interrelationship is sought between  $\varepsilon$ ,  $V_0$ , and the reported  $T_q$  vs  $\bar{M}_n^{-1}$  data, nevertheless, reservations persist that either the variability in the literature (either because of different observers, instrumentation, or sample preparation 9.20.22) or the assumptions inherent in theory (e.g. the failure to account for nearest neighbour interactions<sup>1,23</sup>) may preclude a conclusive interpretation of results such as those addressed in this report.

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 $<sup>^</sup>b$   $K_g$  in K x  $10^{-5}$ 

 $c \in \inf_{\epsilon \text{ in kcal (mol segment)}} -1$ 

d Values for  $T_q^{\infty}$  (K)

<sup>\*</sup> Note error on p 123 of ref 22.  $\varepsilon$  on line 37 should read 0.99 and not

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# Poly( $\beta$ -ketothioethers) by polycondensation of bis- $\beta$ -dialkylaminoketones with bis-thiols

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 $\beta$ -Dialkylaminoketones (Mannich bases) may react with thiols to give  $\beta$ -ketothioethers. Substitution of an -SR group for a dialkylamino group takes place<sup>1-3</sup>:

 $R^{1}$ —CO—CH— $CH_{2}$ —N  $R^{3}$  + H—S— $R^{5}$   $R^{1}$ —CO—CH— $CH_{2}$ —S— $R^{5}$  + H—N  $R^{4}$ 

We have found that by using bifunctional compounds this reaction can open a synthetic route to a novel class of sulphur-containing polymers, the poly( $\beta$ -ketothioethers):

The aim of this Report is to relate some results obtained using 2,6-bis(dimethylaminomethyl)cyclohexanone (I), and 1,2-ethanedithiol (II), oxydiethanedithiol (III) or 1,3-dimercaptobenzene (IV).

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#### **EXPERIMENTAL**

<sup>1</sup>H-n.m.r. spectra were run in CDCl<sub>3</sub> on a C-60 Jeol Spectrometer, with TMS as internal reference.

Bis-Mannich base I was prepared according to Barrett<sup>4</sup>. It was distilled just prior use, b.p. 90–92/0.4 mm.

Dithiols II and III were commercial products (Merck). They were purified by vacuum distillation just before use.

Dithiol IV (Fluka) was used without further purification.

Poly(β-ketothioether) V (Table 1) was prepared by dissolving I (10.0056 g, 47.13 mmol) and II (4.4395 g, 47.13 mmol) in a 3:1 dioxane/ethanol mixture (43 ml) and maintaining the reaction temperature at 60°C by means of a thermostatic alloy-controlled bath. A slow stream of nitrogen was passed through the reaction mixture, in order to prevent oxidation of the dithiol, and to carry out all dimethylamine which was formed during the reaction. The effluent gases were bubbled into aqueous 0.1 N  $\rm H_2SO_4$  (700 ml), and the progress of the reaction was followed by measuring the amount of dimethylamine absorbed by the acid titrimetrically.

After 40 h, 93.7% of the theoretical amount of dimethylamine had been evolved. The reaction was then brought to completion by raising the temperature to 90°C for 3 h. After this time, the solvents were eliminated by distillation under nitrogen, and the product was purified by dissolving the gummy residue in dichloromethane, and reprecipitating with an excess of light petroleum. The polymer was then recovered by filtration, and dried at room